

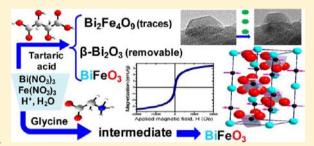


Easy Synthesis of High-Purity BiFeO₃ Nanoparticles: New Insights Derived from the Structural, Optical, and Magnetic Characterization

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Supporting Information

ABSTRACT: Synthesis of high-purity BiFeO₃ is very important for practical applications. This task has been very challenging for the scientific community because nonstoichiometric $Bi_xFe_yO_z$ species typically appear as byproducts in most of the synthesis routes. In the present work, we outline the synthesis of BiFeO₃ nanostructures by a combustion reaction, employing tartaric acid or glycine as promoter. When glycine is used, a porous BiFeO₃ network composed of tightly assembled and sintered nanocrystallites is obtained. The origin of high purity BiFeO₃ nanomaterial as well as the formation of other byproducts is explained on the basis of



metal—ligand interactions. Structural, morphological, and optical analysis of the intermediate that preceded the formation of porous BiFeO₃ structures was accomplished. The thorough characterization of BiFeO₃ nanoparticles (NPs) included powder X-ray diffraction (XRD); scanning electron microscopy (SEM) and high resolution transmission electron microscopy (HRTEM); thermogravimetric analysis (TGA); UV—vis electronic absorption (diffuse reflectance mode), Raman scattering, Mössbauer, and electron paramagnetic resonance (EPR) spectroscopies; and vibrating sample magnetometry (VSM). The byproducts like β -Bi₂O₃ and 5 nm Bi₂Fe₄O₉ NPs were obtained when tartaric acid was the promoter. However, no such byproducts were formed using glycine in the synthesis process. The average sizes of the crystallites for BiFeO₃ were 26 and 23 nm, for tartaric acid and glycine promoters, respectively. Two band gap energies, 2.27 and 1.66 eV, were found for BiFeO₃ synthesized with tartaric acid, obtained from Tauc's plots. A remarkable selective enhancement in the intensity of the BiFeO₃ A₁ mode, as a consequence of the resonance Raman effect, was observed and discussed for the first time in this work. For glycine-promoted BiFeO₃ nanostructures, the measured magnetization (*M*) value at 20 000 Oe (0.64 emu g⁻¹) was ~5 times lower than that obtained using tartaric acid. The difference between the *M* values has been associated with the different morphologies of the BiFeO₃ nanostructures.

1. INTRODUCTION

Bismuth ferrite (BiFeO₃), in its bulk form, is an antiferromagnetic, ferroelectric, and multiferroic material. It has an antiferromagnetic Néel temperature of 643 K and a ferroelectric Curie temperature of 1103 K.¹ Due to these exceptional physical features, great interest has been raised for this material in several fields of science and technology. In the form of thin films, this compound has many potential device applications exploiting its photovoltaic, 2,3 ferroelectric, piezoelectric, or magnetoelectric capabilities, as well as in spintronics. 4 Probably, that is why the thin film has been the most studied and applied form of BiFeO₃. For example, in 2007 Fujitsu introduced a new ferroelectric random access memory (FRAM) for a 90-nm technology node, based on a 5% Mn-doped BiFeO₃ thin film.⁵ For the aforementioned applications, the purity of BiFeO₃ was very critical because the reaction byproducts could modify the magnetic properties.6

Nanoparticulated BiFeO₃ constitutes a very important material that has had less attention paid to it by the scientific

community. Being a semiconductor with the absorption edge in the visible region, BiFeO₃ NPs have been successfully tested in the degradation of rhodamine B,⁷ methyl orange,^{8,9} methylene blue,¹⁰ and bisphenol A.¹¹ Also, BiFeO₃ NPs have been considered for the production of hydrogen through solar water splitting.¹² For this kind of surface reaction on BiFeO₃ NPs, it is important to quantify the remaining NO₃¹⁻ or other anions that could be adsorbed onto the NP surface, thus changing its surface reactivity.¹¹ Furthermore, the impurities or amorphous phases could modify the electron—hole recombination kinetics, affecting the photocatalytic efficiency of BiFeO₃. Therefore, developing methods for synthesizing high purity BiFeO₃ NPs is highly desirable.

In a review by Silva et al., ¹³ a list of secondary phases (Bi₂₅Fe₂O₃₉, Bi₂₅FeO₄₀, Bi₂Fe₄O₉, Bi₄₆Fe₂O₇₂, Bi₃₆Fe₂₄O₅₇), which are common byproducts in the synthesis of BiFeO₃ or

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doped BiFeO $_3$ NPs, were put together after consulting at least 20 papers; these were obtained by a conventional solid state reaction or sol—gel methods. Other additional publications, not included in that review, have also reported the formation of some of those secondary phases, $^{7,14-16}$ an amorphous shell of 1–2 nm thickness around the BiFeO $_3$ NPs, 17 Bi $_2$ O $_3$, 7,18,19 and 5% of an unidentified impurity phase. 13 Lebeugle et al. have pointed out that the Bi $_{25}$ Fe $_2$ O $_{39}$ impurity is responsible for the weak ferromagnetism at low temperature in polycrystalline BiFeO $_3$. 20 However, we found two publications in which few samples synthesized under specific conditions have XRD patterns of pure BiFeO $_3$.

Until now, the most comprehensive characterization for the BiFeO₃ NPs has been documented by Park et al.^{23,24} They have demonstrated the dependence of magnetic properties on the size of NPs. However, there were no band gap estimations, nor TGA experiments, to find out residues from the reagents which can modify the catalytic, magnetic, or optical properties of the NPs. While EPR measurements were not performed, their reported Raman measurements were made with just one laser source of 514.5 nm.²⁴ In the literature, there are some reports which show the Raman scattering spectrum of BiFeO₃, ²⁵⁻³⁴ but only two of them^{26,34} measured the two-phonon overtones using resonance Raman scattering; none of them reported the band gap value either of bulk single crystal or of thin films. When the frequency of the laser beam is close to the frequency of an electronic transition, scattering enhancements have been observed.³⁵ This phenomenon is called resonance Raman scattering, where intense overtones are allowed. The increase in intensity from resonance enhancement can be understood by studying the Kramer-Heisenberg-Dirac (KHD) expression for the polarizability, comprehensively discussed in ref 35. So far, for BiFeO₃, there are no reports about the relationship of overtones and electronic transitions induced by laser sources.

Da Silva et al.³⁶ have reported the preparation of BiFeO₃ NPs via mechanochemical method employing an α-Fe₂O₃/Bi₂O₃ mixture, but the size distribution of their NPs was relatively broad (ranged from about 5 to 40 nm), and an amorphous phase was detected by XRD. However, the maximum magnetic susceptibility at 20 K in the zero field cooling (ZFC) curve, reported by Da Silva et al., differs from that reported at 50 K by Prado et al.³⁷ and that obtained at 85 K by Park et al.^{23,24} At the present, the reported band gap energy values for BiFeO3 NPs and thin films are contradictory with regard to the expected qualitative dependence on the material dimensions (see Table S1). $^{8-10,25,38-49}$ It should be expected that the smaller the NPs the higher the band gap energy value, but the reported values do not show this trend. In the particular case of BiFeO₃ NPs, there is no agreement in theoretical considerations whether the electronic transitions are direct or indirect, see Table S1. This fact might introduce discrepancies among the experimentally reported band gap values.

The main objectives of the present research are to synthesize, purify, and characterize BiFeO₃ NPs, as well as the impurities (even at trace levels), with a wide variety of spectroscopic, structural, morphological, thermal, and magnetic techniques. We developed an easy method to obtain high purity BiFeO₃ nanocrystals at relatively low temperatures. This process is straightforward, simple, energy-saving, and cost-effective. An integral discussion about the role of glycinate ion as ligand, which has the same affinity constant for Fe³⁺ and Bi³⁺ (log K = 10)^{50,51} and slightly lower for H⁺ (log K = 9.57),⁵¹ was accomplished. After a thorough characterization of the BiFeO₃

NPs by powder XRD, SEM, HRTEM, TGA, VSM, and UV—vis electronic absorption, Raman scattering, Mössbauer, and EPR spectroscopies, we obtained very interesting experimental results and offer consistent interpretations of the physicochemical properties of these ferrite NPs.

2. EXPERIMENTAL SECTION

The reagents used in the synthesis were bismuth nitrate pentahydrate (Bi(NO₃)₃·SH₂O, Aldrich, > 98%), iron nitrate nonahydrate (Fe(NO₃)₃·9H₂O, J.T. Baker, 99.6%), L-(+)-tartaric acid ($C_4H_6O_6$, Alfa Aesar, 99%), glycine (H_2NCH_2COOH , Aldrich, 99%), diluted nitric acid (HNO₃, J.T. Baker, 64%), and glacial acetic acid (CH₃COOH, J.T. Baker, 99.9%).

A typical preparation process used to obtain the precursors for the synthesis of BiFeO $_3$ NPs is the following: 0.004 mol of Bi(NO $_3$) $_3$: $5H_2O$, 0.004 mol of Fe(NO $_3$) $_3$:9 H_2O , and 0.008 mol of L-(+)-tartaric acid or glycine were dissolved in a mixture of 40 mL of deionized water and 2.8 mL of HNO $_3$ at 64%, under magnetic stirring at room temperature. In each reaction, only one organic species (either tartaric acid or glycine) was used to complex Bi $^{3+}$ and Fe $^{3+}$ cations in aqueous solution. The mixture was heated until all solvent was evaporated, and a combustion reaction took place. In the reaction, the organic species was the reducing agent, and the NO $_3$ ¹⁻ was the oxidizing agent. The brown powder obtained using tartaric acid was heated in a furnace at 500 °C in air for 1 h, and then it was washed with water followed by CH $_3$ COOH at 65 °C to remove the Bi $_2$ O $_3$. The powder obtained using glycine was annealed in air at 350 °C for 1 and then at 500 °C for 1 h in addition.

The measurements of the electronic absorption spectra using the diffuse reflectance technique were made in a UV–vis spectrophotometer (Ocean Optics CHEM-2000), equipped with a double-way optic fiber, coupled to a PC. The powder XRD patterns were collected in a Siemens DS000 diffractometer using a Cu source (K α , λ = 1.5406 Å). The HRTEM observation was carried out in a Philips Tecnai F30 microscope operated at 300 KeV and equipped with a Gatan Image Filter to perform the EELS study. HRTEM images were processed using the Digital Micrograph 1.2 software. The SEM images were obtained in a field emission scanning electron microscope, JSM-7600F, from JEOL, with a secondary electron detector. The Raman scattering spectra were measured in a Horiba spectrometer, equipped with 3 laser sources (λ = 532, 633, and 785 nm), and a Synapse CCD detector with thermoelectric cooling at -75 °C.

The magnetic hysteresis loops (M-H) were measured in a vibrating sample magnetometer (VSM-VersaLab, 30000 Oe), at room temperature. The magnetic susceptibility at different temperatures was measured in a superconducting susceptometer MPMS-5s, Quantum Design. The thermograms were recorded on a TGA Q5000 V3.10 Build 258 equipment, brand TA, in modulated mode, using a heating ramp of 5 °C/min. The Mössbauer spectrometer (Wissel-Electronik) with a 57 Co source, embedded in a Rh matrix, with activity of 925 MBq (25 mCi), was operated at 300 and 77 K. The spectra fittings were made with the NORMOS software. EPR spectra were recorded on a Bruker equipment (Elexsys E 500).

3. RESULTS AND DISCUSSION

3.1. X-ray Powder Diffraction. *a.* $BiFeO_3$ NPs Synthesized Using Tartaric Acid as Ligand. Typical X-ray powder diffraction (XRD) patterns of the BiFeO₃ obtained using tartaric acid before and after its purification are shown, Figure 1. The peaks in Figure 1a belong to rhombohedral BiFeO₃ (PDF 861518), and the main byproducts were β -Bi₂O₃ (PDF 781793)⁵³ as indicated with the blue arrows, and the α -Bi₂O₃ phase (PDF 761730) as indicated with the asterisk.

Figure 1b shows an XRD pattern of the BiFeO₃ powder, taken after washing with CH₃COOH to remove the bismuth oxide impurities. Usually these oxides were removed with HNO₃, ¹⁸ a strong acid, which could dissolve BiFeO₃. For the

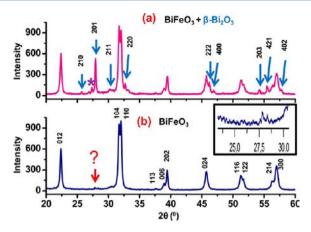


Figure 1. X-ray powder diffraction patterns of BiFeO₃ synthesized using tartaric acid as ligand (a) before and (b) after the treatment with CH₃COOH to remove the byproducts: β -Bi₂O₃ and α -Bi₂O₃. The inset is a zoom-in around the peak at $2\theta = 27.90^{\circ}$.

first time, a weak acid, CH3COOH, was used for purifying purposes. The average crystallite size of BiFeO₃ nanostructures was 26 nm, determined by means of the Scherrer equation, using the (012) and (024) planes only. The relatively small peak located at $2\theta = 27.91^{\circ}$, indicated with the question mark and shown in the inset, does not belong to BiFeO₃. This peak matches with the most intense reflection of Bi24Fe2O39, located at $2\theta = 27.97^{\circ}$ (PDF 420201). The signal at $2\theta = 27.97^{\circ}$ is also close to the most intense peak of Bi₂Fe₄O₉ (PDF 250090), a common byproduct, which is located at $2\theta = 28.20^{\circ}$ (PDF 250090). Recently, Kothai et al. 54 reported the Bi₂₄FeO₃₉ impurity during the synthesis of BiFeO3 NPs based on a small peak located at $2\theta = 27.9^{\circ}$. Considering that an amount of bismuth was consumed for the β -Bi₂O₃ formation, we propose Bi₂Fe₄O₉ as the most probable byproduct. In section 3.3 and Figure 6, there is a further discussion about the question mark

We suggest the following pathway to explain the formation of the bismuth oxide phases as byproducts. In Bi(III) acidic aqueous solutions, intramolecular polycondensation of the hydrated ion takes place; leading to polynuclear cationic species such as $[\mathrm{Bi}_6\mathrm{O}_4(\mathrm{OH})_4]^{6+}$, $[\mathrm{Bi}_6\mathrm{O}_5(\mathrm{OH})_3]^{5+}$, and $[\mathrm{Bi}_9(\mu_3\text{-}\mathrm{O})_8(\mu_3\text{-}\mathrm{OH})_6]^{5+}$. These polynuclear species can form clusters and then, under thermal treatment, decompose into the $\alpha\text{-}$ and $\beta\text{-}\mathrm{Bi}_2\mathrm{O}_3$ phases. This is the first time that the formation of the byproducts $\alpha\text{-}$ and $\beta\text{-}\mathrm{Bi}_2\mathrm{O}_3$ could be associated with these polynuclear species in the synthesis of BiFeO_3 NPs.

The anion of the L-(+)-tartaric acid and Bi^{3^+} form polymeric chains with a nine-coordinate bismuth center $(\mathrm{BiO_9})^{.51,61}$. However, the tartrate anion does not have the same affinity for Bi^{3^+} and Fe^{3^+} . S1,62,63 Consequently, the two cations do not get homogeneously distributed in the polymeric chains in the precursor solutions. This could be a new possible explanation for the formation of $\mathrm{Bi_2O_3}$ and $\mathrm{Bi_2Fe_4O_9}$ byproducts, or a different $\mathrm{Bi/Fe}$ atomic ratio from 1:1 in the $\mathrm{BiFeO_3}$ NPs, obtained during the combustion reaction.

b. BiFeO₃ NPs Synthesized Using Glycine as Ligand. Figure 2 shows the XRD patterns of BiFeO₃ NPs at different stages of the heating process. In Figure 2a, several peaks do not belong to the BiFeO₃ phase, and the most intense peak is located at $2\theta = 28.17^{\circ}$. Formation of an amorphous phase can also be observed from the curvature of the baseline. Although the most intense peak could not be directly assigned to any specific

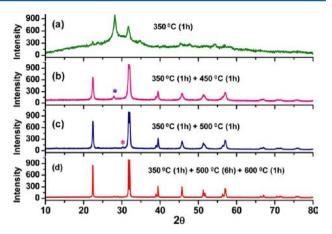


Figure 2. XRD patterns of (a) precursor (intermediate) and (b-d) the BiFeO₃ NPs, obtained using glycine as ligand, annealed at different temperatures and times. The asterisk in part c corresponds to the byproduct $(BiO)_2CO_3$.

species, the Raman spectrum of this material only shows signals corresponding to BiFeO₃ (see Figure S1) suggesting that the sample is mainly constituted by that phase and probably an oxygen deficient bismuth ferrite. The presence of oxygen deficient bismuth ferrite could be due to the formation of combustion gases (NO, NO₂, CO, CO₂, and H₂O) that take up the available oxygen, and the capability of CO₂ and H₂O to remove molecular oxygen. SEM analysis revealed the generation of a porous framework of nanocrystals smaller than 10 nm in this intermediate (see Figure S2a). It is noteworthy to mention that the crystallites forming this framework are preferentially aligned in the porous assemble (see Figure S2c).

From the ATR-FTIR spectrum in Figure S4, the bands at 846 and 1379 cm⁻¹ confirm the presence of (BiO)₂CO₃, and the remaining bands agree with the NO₃¹⁻ group coordinated to Bi³⁺.⁶⁴ After heating this sample for one additional hour at 450 °C (Figure 2b), all the peaks correspond to the BiFeO₃ phase, except that indicated with the asterisk at $2\theta = 27.96^{\circ}$.

When the sample was heated at 350 °C for an hour and then at 500 °C for an additional hour (Figure 2c), all the prominent peaks that appeared in its XRD pattern correspond to BiFeO₃, except the weak peak appearing at $2\theta = 30.26^{\circ}$ marked by the asterisk in Figure 2c, which was associated to the (BiO)2CO3 phase, as has been reported previously by Taylor et al.⁶⁵ The presence of (BiO)₂CO₃ was also confirmed by its ATR-FTIR spectrum (see Figure S4). The average crystallite size of the BiFeO₃ (Figure 2c) was about 23 nm, determined by the Scherrer equation, only using the (012) and (024) planes. The α - and β -Bi₂O₃ byproducts were not detected by XRD, so the treatment with CH3COOH was not necessary. This is the main difference observed between the use of glycine or tartaric acid for the synthesis. The reasons for this pathway differentiation might be the following: (a) There is a homogeneous distribution of the Fe(III)-glycine and Bi(III)-glycine complexes due to the similar affinity constants that the glycinate anion has for the cations (Fe³⁺, Bi³⁺, and H⁺).^{50,51} (b) The glycinate anion competes with NO₃¹⁻ and interrupts the formation of the 3D networks of bismuth oxide clusters, like $[Bi_6O_4(OH)_4(NO_3)_6(H_2O)] \cdot H_2O$. The latter forms bismuth oxide when decomposed under annealing. (c) The pK_a of the carboxylic group of glycinate ion (2.34) is lower than that of tartaric acid (2.95 and 4.25).

In order to evaluate the thermal stability of the BiFeO₃ NPs, they were heated as indicated in Figure 2d. The heating process resulted in a size increase to about 67 nm (calculated using the Scherrer equation). The peak that corresponds to the (BiO)₂CO₃ phase was no longer detected after the annealing process as it decomposes between 350 and 650 °C.⁶⁵ The high purity of the obtained BiFeO₃ powder observed by XRD was corroborated by TGA analysis; the results are presented and discussed in Figures S5 and S6.

3.2. Scanning Electron Microscopy (SEM). *a.* BiFeO₃ NPs Synthesized Using Tartaric Acid as Ligand. For further characterization, the BiFeO₃ NPs were analyzed by SEM, and a typical image is presented in Figure 3. It is clear that the NPs

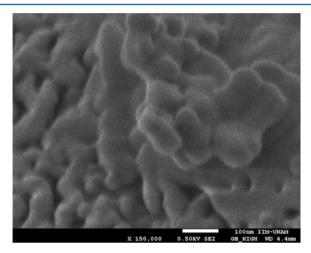


Figure 3. Typical SEM image of the aggregated $BiFeO_3$ NPs, prepared using tartaric acid. The scale bar is of 100 nm.

are aggregated and have sizes smaller than 50 nm, which agree with their XRD results. The Bi/Fe ratio was determined by the SEM–EDS technique. These ratios were measured in two different zones of an area of about 1 μ m² each, where 0.91 and 1.04 should be 1 in the ideal case; these values are acceptable within the accuracy of the technique and comparable with previous works. For example, for the BiFeO₃ films deposited at low oxygen pressures, the ratio was calculated to be about 1.2:1 using XPS measurements.⁴⁹

b. BiFeO₃ NPs Synthesized Using Glycine as Ligand. In the images of BiFeO₃ NPs analyzed by SEM, it looks like a highly porous network instead of simple nanoparticle aggregates, as can be seen in Figure 4. The pore sizes ranged from several micrometers to submicrometer dimensions, and these holes have irregular shapes. The pores are created due to the fast release of combustion gases. This structure is derived from the previously formed intermediate (Figure S2a,b). In the literature we can find that, during the combustion reaction between glycine and metal nitrates, the temperature inside the flame is increased to a maximum between 1100 and 1450 °C, 52,66 rising from 150 to 1285 °C in less than 4 s;⁵² consequently, several gases are abruptly released (NO, NO₂, CO, CO₂, NH₃, and H_2O), $S_2,67$ and products with microporous structures are formed.^{7,68–70} Hwang et al. have reported a combustion reaction between metal (Ni, Zn, and Fe) nitrates and glycine studied by differential scanning calorimetry (DSC) in which the curve showed a sharp exothermic peak at ~178 °C with an area of -520 cal/g.^{52} This energy can be taken by the released gases to create those microporous structures observed in refs 7, 68,

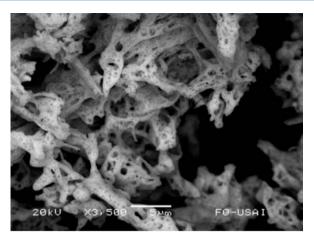


Figure 4. SEM image of the BiFeO₃ porous network prepared using glycine in the synthesis methodology. The scale bar is of 5000 nm.

69, and 70. It can be inferred from the experiments in ref 52 that 64% of the sample weight evolves as gases in about 0.25 s, while the microporous structure is formed.

The theoretical values for the atomic percentage of Bi, Fe, and O in BiFeO₃ are 20, 20, and 60, respectively. These values are close to those obtained experimentally by SEM-EDS (19 \pm 2%, 20 \pm 4%, and 60 \pm 4%).

3.3. High Resolution Transmission Electron Microscopy (HRTEM). *a.* $BiFeO_3$ NPs Synthesized Using Tartaric Acid as Ligand. An HRTEM image of one BiFeO₃ NP is shown in Figure 5a. This NP has an ellipsoidal shape and approximate dimensions of $14 \times 19 \text{ nm}^2$. The interplanar spacing observed in this micrograph is 3.92 Å and corresponds to the $\{012\}$ BiFeO₃ family planes, that is, the first peak shown in Figure 1b. Figure 5b corresponds to the fast Fourier transform (FFT) of the selected area, in which the interplanar distances of the BiFeO₃ were identified. The reciprocal lattice simulated by the CaRIne software, viewed from the [14-2] zone axis, shown in Figure 5c, agrees well with the FFT of Figure 5b.

The only byproduct found in some HRTEM images was tiny NPs of about 5 nm size, as shown in Figure 6a. The FFT analysis of the zone marked by the red square in Figure 6a is presented in Figure 6b. The interplanar distance found was 3.2 Å. This distance matches with the interplanar distance associated to the very small peak indicated with the arrow in the diffractogram of Figure 1b, which means that these tiny NPs are possibly $\mathrm{Bi}_2\mathrm{Fe}_4\mathrm{O}_9$.

The iron-rich $Bi_2Fe_4O_9$ impurity was possibly formed since a fraction of bismuth atoms were consumed in the formation of the β -Bi₂O₃ byproduct, keeping a small amount of Fe(III) in the reaction media. The Bi₂Fe₄O₉ byproduct had already been reported; ^{7,54,72,73} however, this is the first time that its size was determined, which is very important because at low temperatures the saturation magnetization of $Bi_2Fe_4O_9$ NPs increases as the size decreases. ⁷⁴ Another relevant result obtained was that the tiny $Bi_2Fe_4O_9$ particles underwent a rearrangement of the crystalline structure through quasimelt states (structural instability) induced by the electron beam from the microscope, ^{75–77} as can be seen in the video in the Supporting Information. Da Silva et al. ³⁶ prepared BiFeO₃ NPs via mechanochemical processing of an α -Fe₂O₃/Bi₂O₃ mixture, and they highlighted the formation of small NPs with interplanar distance of 3.2 Å that were formed by a rapid

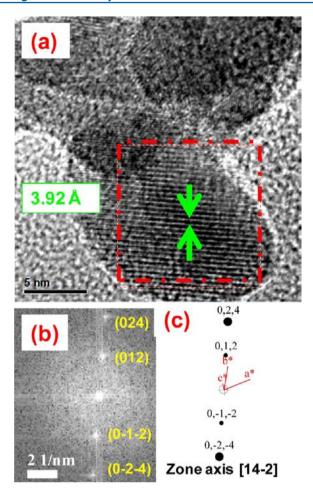


Figure 5. (a) Typical HRTEM image of one BiFeO₃ NP. (b) Fast Fourier transform of the squared area marked in part a. (c) Electron diffraction pattern simulated by CaRIne v. 3.1 for BiFeO₃.

crystallization reaction upon the electron irradiation during the course of HRTEM investigation. We suggest that what they saw in their HRTEM images was just structural instability but not formation of small NPs. A careful observation of the micrographs in ref 36 revealed the presence of those NPs even since t=0 of electron-beam irradiation. In our case, the NPs' shape and facets notably changed as the lattice was rearranged by the electron beam interaction, including the so-called quasimelt states.

b. BiFeO₃ Porous Network Synthesized Using Glycine as Ligand. The HRTEM images of the BiFeO₃ NPs, obtained using the glycine methodology, are shown in Figure 7a, and at least in that image, their sizes are smaller than 60 nm and are aggregated. It can be seen that the BiFeO₃ NPs were tightly assembled and sintered giving support to the porous network presented in Figure 4. Hereafter, this BiFeO₃ NPs porous network will be called BiFeO₃—NPs—PNt. A closer view of one BiFeO₃ NP allowed us to determine the interplanar distances and confirm that it corresponds to the {104} family of planes, as can be seen in Figure 7b. Surrounding this BiFeO₃ NP, there are some unidentified small NPs, two of them indicated with the blue arrows.

3.4. Comparison of the Electronic Absorption Spectra of the BiFeO₃ NPs Prepared Using Glycine and Tartaric Acid as Ligands. The solid state UV—vis electronic absorption spectra of BiFeO₃ NPs at different annealing stages

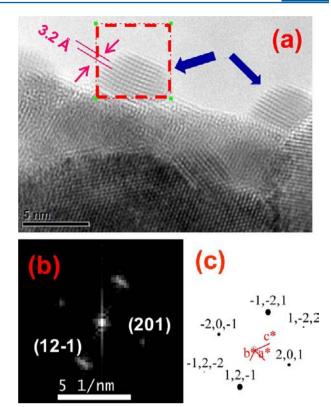
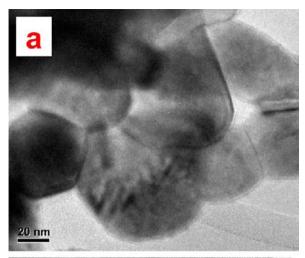


Figure 6. (a) HRTEM image of a BiFeO $_3$ sample where two NPs of the byproduct Bi $_2$ Fe $_4$ O $_9$ were identified. (b) Fast Fourier transform of the zone indicated with the red square in part a. (c) Electron diffraction pattern simulated by CaRIne v. 3.1 for Bi $_2$ Fe $_4$ O $_9$ using the zone axis [-234].

are shown in Figure 8. The sample with the portion of amorphous phase (shown in Figure 2a) has the less steep absorption edge in its absorption spectrum between 500 and 600 nm. This attribute is probably due to a variety of electronic environments around the iron(III) and bismuth(III) cations, possibly created by the remaining nitrates and carbonates (see Figure S3). Additionally, several other factors, including the amorphous phase, the smaller BiFeO₃ nanostructures, higher lattice strain, and the presence of other secondary phases, contribute to the resulting absorption profile. When the temperature is increased, the amount of nitrates decreases, and the percentage and the size of BiFeO3 NPs increase while the lattice strain decreases, and hence make the absorption edge steeper. The nanostructures obtained at the longest annealing time are of largest in size, i.e., 67 nm, and have the most abrupt absorption edge. This can be explained since these nanostructures are the most crystalline and purest; they must have the smallest specific surface area and the highest crystalline order at the interparticle contact. This means less absorption contribution from the tails density of states. Taking into account that the heating treatment was performed in air atmosphere, these largest crystallites should have fewer oxygen vacancies.

The spectrum corresponding to the BiFeO $_3$ NPs prepared using tartaric acid as ligand (curve a) has a profile very similar to that obtained using glycine (curve d), see Figure 8. This is in accordance with the very similar crystallite size of both samples (26 and 23 nm, respectively).

From the linear fit in the Tauc's plot, restricted to the energy range 2.17-2.36 eV (525-570 nm), a direct band gap of 2.27 eV is found for the BiFeO $_3$ nanostructures of 23 nm crystallite



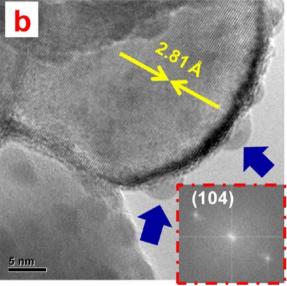


Figure 7. HRTEM images of (a) BiFeO₃–NPs–PNt and (b) unidentified small NPs surrounding a BiFeO₃ NP (two of them are indicated by the arrows); the inset in part b corresponds to the FFT. The corresponding XRD pattern of this sample is shown in Figure 2c.

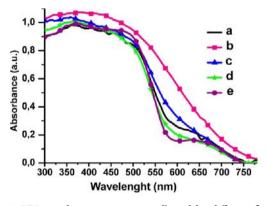


Figure 8. UV—vis absorption spectra collected by diffuse reflectance technique of BiFeO₃ NPs obtained using tartaric acid (a) and glycine (b, c, d, e) as ligands. The times and temperatures of heating were the following: (a) $500 \,^{\circ}$ C, 1 h; (b) $350 \,^{\circ}$ C, 1 h; (c) $350 \,^{\circ}$ C, 1 h, + $450 \,^{\circ}$ C, 1 h; (d) $350 \,^{\circ}$ C, 1 h, + $500 \,^{\circ}$ C, 1 h; (e) $350 \,^{\circ}$ C, 1 h, + $500 \,^{\circ}$ C, 6 h, + $600 \,^{\circ}$ C, 1 h.

size (see Figure 9a). This direct band gap value is in-between those reported for BiFeO₃ nanostructures (1.8-2.5

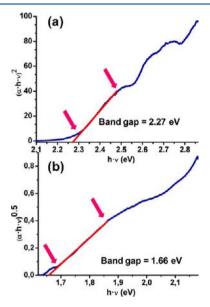


Figure 9. Tauc's plots to determine the band gap of the BiFeO₃ NPs, prepared using tartaric acid as ligand agent: (a) direct electron transition and (b) indirect electron transition.

eV) $^{9,25,78-80}$ that, at the same time, are red-shifted with respect to the bulk value $(2.66-2.81 \text{ eV}).^{48,81-83}$ This effect has been explained on the basis of the fact that the typical electric dipole-allowed charge transfer excitations of bulk BiFeO₃ at 3.2 and 4.5 eV split into multiplets in nanoparticles with the net effect of moving the edge of the charge transfer band to lower energies. That is a consequence of the strained nature of the nanostructures and the symmetry breaking in high surface to volume ratio systems. The red shift of the absorption spectrum in BiFeO₃ has also been observed in other lattice-strained, doped, or surface-rich structures and BiFeO₃ nanowires.

Additionally, an indirect interband transition was detected at 1.66 eV in the same BiFeO₃ nanostructures of 23 nm crystallite size. It was derived from the Tauc's plot fit in the long tail-like part of the absorption spectrum, at wavelengths longer than 580 nm. The presence of both indirect and direct band gaps, in different energy ranges, has been observed by other authors in BiFeO₃ thin films. 48,88 In these cases the indirect transitions are located between 1.3 and 1.85 eV. 48,88 These reports are qualitatively in accordance with some theoretical studies where the existence of an indirect and a direct gap are predicted in the electronic structure of BiFeO₃.4,72 The appearance of an indirect band gap has been also theoretically predicted in highly strained BiFeO₃. ⁸⁵ In these reports, ^{4,72,85} the indirect gap is of smaller energy than the direct one. ⁸⁹ According to Wang et al.,89 the top of the valence band is composed by strong hybridization among O 2p, Bi 6p, and Fe 3d states, and the bottom of the conduction band is leading by Fe 3d states. Very similar band gap values, for direct and indirect transitions, were determined for BiFeO₃-NPs-PNt, see Figure S8.

3.5. Raman Scattering Spectroscopy. *a. BiFeO*₃ *NPs Synthesized Using Tartaric Acid as Ligand.* BiFeO₃ belongs to the R3c space group and the $C_{3\nu}$ point group. The irreducible representations obtained from group theory are the following: $4A_1(z, x^2 + y^2, z^2) + 5A_2(-) + 9E(x, y, x^2 - y^2, xy, xz, yz).$

However, the A₂ Raman modes are nonactive. In Figure 10 the measured Raman scattering spectra of the BiFeO₃ NPs, using

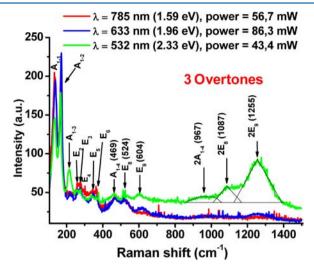


Figure 10. Raman spectra of the 26 nm ${\rm BiFeO_3}$ NPs, prepared using the tartaric acid methodology, employing three different lasers as an excitation source.

three laser sources with different energy values, are presented. All bands in the spectra belong to the BiFeO₃ NPs.²⁵⁻³⁴ Using the laser source of 532 nm (2.33 eV), which is just 0.06 eV

higher than the band gap value of the NPs (2.27 eV, as presented in Figure 9a), three overtones were observed and that at 1255 cm $^{-1}$ is about 10 times more intense than the corresponding fundamental mode (E₉). The intensity of the overtone at 1255 cm $^{-1}$, using the laser source of 633 nm (1.96 eV), was very low because it only has enough energy for indirect electron transitions, which are about 1000 times less intense than the direct electronic transitions. When the laser source of 785 nm (1.59 eV) was used, the overtones were absent. These results confirm that electronic transitions take place under excitation with 532 nm laser radiation, as expected from the diffuse reflectance measurement, allowing the resonance effect.

Another remarkable finding was a notable intensity enhancement of the fundamental totally symmetric mode at 216 cm⁻¹ (see Figure 11a), labeled as A₁₋₃, when the laser source of 532 nm was used. It is important to mention that, in the resonance Raman phenomenon, the intensity enhancements observed in some vibrational modes are exclusively related to the atoms taking part in the electronic transition. ³⁵ According to Porporati et al., ⁹¹ Bi atoms only participate in the low wavenumber modes up to 167 cm⁻¹, while the oxygen motion strongly dominates in modes above 262 cm⁻¹ and Fe atoms are mainly involved in modes between 152 and 261 cm⁻¹. Moreover, first-principles calculation within generalized gradient approximation made by Liu et al. ⁹² and Wang et al. ⁸⁹ concluded that the density of states at the lower conduction bands are conformed

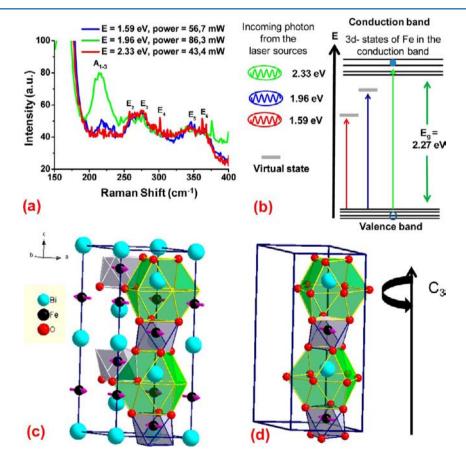


Figure 11. (a) Intensity enhancement of the A_1 Raman mode at 216 cm⁻¹ for BiFeO₃ using three laser sources of different energy (E), (b) sketch of the relationship between the energy of the laser sources used in the Raman measurements and the energy of the band gap, (c) sketch of the BiFeO₃ structure in the hexagonal unit cell using the atomic parameters obtained by Kubel et al., 93 and (d) representative polyhedra around the Fe³⁺ and Bi³⁺ cations.

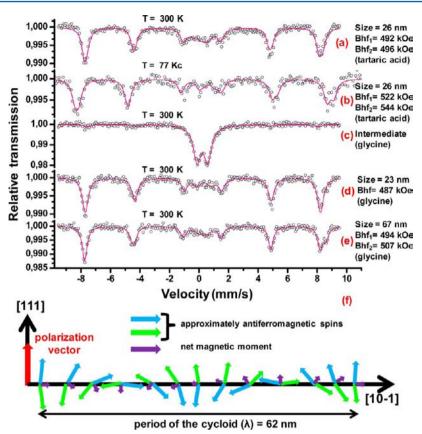


Figure 12. Mössbauer spectra of BiFeO₃ NPs of different sizes prepared using tartaric acid (a, b) and glycine (c, d, e) in the synthesis. Spectrum c corresponds to the intermediate. The spectra a, c, d, and e were measured at 300 K, while spectrum b was recorded at 77 K. (f) Schematic representation of the spin cycloid inspired from ref 4.

mostly by Fe-d states (see Figure 11b). These facts could mean that when the laser's energy is sufficient to promote an electronic transition from the valence to the lower conduction bands in $BiFeO_3$, a perturbation occurs in the electronic density of the 3d orbitals of the Fe atoms, and consequently, changes in the polarizabilty of iron atoms are manifested as an enhancement in the intensity of only the A_1 mode at 216 cm⁻¹.

The *R3c* symmetry in BiFeO₃ permits the development of spontaneous polarization (Ps) along the pseudocubic [111] direction in the cell, due to the displacement of Bi³⁺ and Fe³⁺ cations from their centrosymmetric positions along the 3-fold rotation axis. Also, BiFeO₃ can be viewed as a hexagonal cell, instead of the rhombohedral unit cell; for example, using the atomic parameter obtained by Kubel, the sketch in Figure 11c was obtained. The polyhedron around the Fe³⁺ shares an equilateral triangle above and below with the polyhedron around Bi³⁺, as can be seen in Figure 11d (some oxygen atoms of the neighbor hexagonal cells were added to see the atomic environments around the bismuth and iron cations more clearly). The bismuth cations are slightly above the center of their polyhedron due to the stereochemical consequences of the bismuth atom electron lone pair.

When an electronic transition takes place, it can induce a potential energy minimum displacement along the normal coordinate between the ground and the excited electronic states. This condition guarantees the non-nullity of the vibrational overlap integral (Franck–Condon factor) in the resonance Raman effect. Small potential energy minimum displacement provokes enhancements only in the fundamental vibrations. This is the case already discussed at 216 cm⁻¹

 (A_{1-3}) wavenumber, mainly involving Fe atoms. As the mentioned displacement increases, overtones are also observed, and their intensities increase relative to the intensity of the fundamental vibration. That means, in our case where overtones are located between 800 and 1500 cm⁻¹, oxygen atoms have the weightiest role in the potential energy minimum displacement when electronic transitions are excited at 2.33 eV. At the same time, the higher the frequency of the fundamental vibration, the higher the intensity for the corresponding overtone; so, potential energy minimum displacement is strongly related to those higher frequency modes.

3.6. Mössbauer Spectroscopy. Figure 12 shows the asymmetric Mössbauer spectra of the BiFeO₃ NPs synthesized using tartaric acid and glycine. It can be seen that the magnetic sextet contribution to the experimental spectra (Figure 12a,b,d,e) has slightly broader resonance lines with a noticeable asymmetry. For further discussion on the possible origin of the spectra asymmetry, see the Mössbauer spectroscopy section in the Supporting Information. The corresponding Mössbauer parameters, after fitting the spectra, are presented in Table S2 in the Supporting Information section. For the NPs obtained using tartaric acid (Figure 12a,b) there are two sextets and two doublets, as reported in previous works.^{23,24} At 300 K, the two sextets, which can only be present in ferro-, ferri-, or antiferromagnetic materials, have slightly different hyperfine magnetic field (Bhf) values (492 \pm 2 and 496 \pm 2 kOe) and account for 79% of the relative area of the spectrum. At 77 K, these two Bhf values increase to 522 and 544 kOe, and the two sextets account for 95% of the relative area.

In a group of papers of ^{57}Fe NMR for BiFeO $_3$ reported by Zalessky et al., $^{95-98}$ the obtained spectra have two peaks of different heights that correspond to two Bhf values measured at 296 K (490.5 and 496 kOe) and at 77 K (538.8 and 546 kOe). The gyromagnetic ratio of ^{57}Fe nuclei ($\gamma/2\pi=0.138$ MHz/kOe) was used to express the NMR frequency values in terms of Bhf values. 95,96 The Mössbauer Bhf values presented in Figure 12a,b agree very well with the ones obtained by ^{57}Fe NMR and, therefore, support the Zalessky's mathematical model.

The two doublets that account for the 21% of the relative area could be related to (a) the subset of BiFeO $_3$ NPs with the smallest size that have superparamagnetic behavior, (b) Fe on the NP's surface coordinated with CH $_3$ COO $^-$, (c) the formation of Bi $_2$ Fe $_4$ O $_9$ byproduct which has 2 doublets at 300 K. ⁹⁹ The documented quadrupole splitting for these two doublets in Bi $_2$ Fe $_4$ O $_9$ are 0.94 and 0.38 mm s $^{-1}$ with isomer shifts of 0.24 and 0.36 mm s $^{-1}$. ^{99,100} These values are similar to our values in Table S2. When the temperature was decreased to 77 K, the relative area also decreased to 5.3% because at this temperature the doublets became two overlapping sextets, as documented by Kostiner et al. ¹⁰⁰ According to Park et al., ^{23,34} in BiFeO $_3$ NPs of size 14 nm, 66% of the relative area fit to two doublets at 300 K, and they become two sextets at 4.2 K. When the size of the NPs increased to 95 \pm 28 nm there were no doublets in the Mössbauer spectrum at 300 K. ²³

It is worthwhile to point out that the intermediate species (Figure 12c) has a low intensity sextet; however, that amount of ferromagnetic iron in the sample is enough to be attracted by a $1.875 \times 0.875 \times 0.393$ in.³ NdBFe magnet. The sextet in the intermediate corresponds to the BiFeO₃ already formed; while the doublet corresponds to high spin ferric iron.

The Mössbauer spectrum of the BiFeO₃–NPs–PNt with 23 nm of crystallite size (Figure 12d) was fitted to just one sextet with a Bhf value of 494 kOe and 88% of the relative area. On the other hand, the spectrum of the BiFeO₃–NPs–PNt of larger crystallite size (67 nm) has broader absorption bands and was fitted to two sextets with Bhf values of 494 and 507 kOe comprising 87% of the relative area (Figure 13e). This difference is probably due to the size, since 67 nm is slightly larger than the period length of 62 nm of the known SSMS for bismuth ferrite while 23 nm is smaller than that period (see Figure 12f).²³

3.7. Vibrating Sample Magnetometry. a. BiFeO₃ NPs Synthesized Using Tartaric Acid as Ligand. The magnetization curve (M-H) for the 26 nm BiFeO₃ NPs is shown in Figure 13a for which a noticeable ferromagnetic response is observed, with a saturation magnetization M_s of 3.4 emu g⁻¹ and a coercive field H_c of 8 Oe. This behavior is consistent with previous reports of nanosized BiFeO₃ particles with mean diameter below 62 nm, whose enhanced ferromagnetism can be ascribed to the effect of breaking periodicity (62 nm) of the spin cycloid structure, characteristic of bulk BiFeO₃. A summary of the magnetic properties is shown in Table S3, including comparative data of other works. According to the results of Park et al.,23 the Hc values increase with the size of BiFeO₃ NPs; so, higher H_c values are explained in terms of the higher size of their BiFeO₃ NPs (compared with our results). Wang et al.²¹ have synthesized BiFeO₃ NPs of 60–90 nm using tartaric acid, and the M value (0.34 emu g⁻¹) at 15 000 Oe was almost 1 order of magnitude lower than the measured value reported in the present work, which is probably due to their larger size.

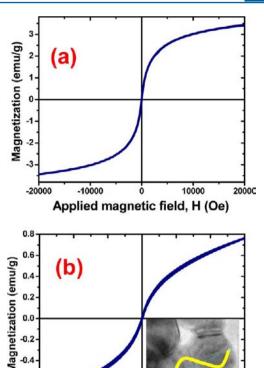


Figure 13. Magnetic hysteresis loop of BiFeO $_3$ NPs using tartaric acid (a) and glycine (b) in the synthesis methodology. Inset in part b: HRTEM micrograph of tightly assembled and sintered BiFeO $_3$ NPs forming the porous network. This illustrates the possibility of sintered neighbor NPs with similar crystalline faces. Superimposed on the micrograph, a sketch of the spin cycloid (62 nm) indicated by the yellow sine function graph.

Applied magnetic field, H (Oe)

20000

-20000 -10000

b. BiFeO₃-NPs-PNt Synthesized Using Glycine as Ligand. For BiFeO₃-NPs-PNt with 23 nm of crystallite size, a clear ferro/antiferromagnetic behavior is manifested in the corresponding M-H curve (Figure 13b). The observed M value at 20 000 Oe (0.64 emu g^{-1}) was ~5 times lower than that obtained using tartaric acid. This result agrees with the fact that no undesirable secondary magnetic phases are present. In addition, the sintered BiFeO₃ nanocrystals forming the porous network (insert in Figure 13b) might promote a partial assembly and continuity of the spin cycloid. This can lead to a dilution effect of the magnetic moment of the powder and, hence, to the significant reduction observed for M (0.64 emu g⁻¹). The strong tendency of preferential alignment during the nanocrystal aggregation has been observed and theoretically supported in different oxides. ^{101,102} This preferential alignment has been explained in terms of interparticle interaction forces and/or external forces. ¹⁰¹ The nature of the ordering physical field-forces can be (i) electric, (ii) magnetic, and/or (iii) higher order multipole moments interaction of the crystalline faces. 101 It is important to highlight that the contact between parallel faces of the hexagonal nanocrystals (insert in Figure 13b) suggests the contact between the faces (10-10), (01-10), or (-1100) of the hexagonal unit cell.⁹³ The mutual alignment of NPs of identical crystal faces is one of the main arguments used in the literature 101 to explain the formation of threedimensional NPs arrays (mesocrystals).

In turn, this preferential alignment tendency of crystallite aggregates is different from what takes place when BiFeO₃ nanopowders are obtained using tartaric acid (see Figure 5a). In such a case, a mixture of BiFeO₃ and Bi₂O₃ is formed in the synthesis (Figure 1a). Then, the direct contact between BiFeO₃ particles can be partially hindered, and the driving forces for the preferential alignment are minimized. Finally, the BiFeO₃ nanocrystallites rinsed with acetic acid do not have strong sintering and exhibit slightly rounded surfaces probably because of chemical etching.

Previous works in the literature that have used EPR for BiFeO₃ characterization always contained Bi₂Fe₄O₉ as impurity in considerable amounts.¹⁰³ To the best of our knowledge, this is the first time that EPR parameters (i.e., resonance field, *g*-factor, and signal width) are reported for high purity BiFeO₃.

CONCLUSIONS

In summary, an integral analysis and discussion of the synthesis pathways and structural, morphological, optical, and magnetic attributes of two sorts of BiFeO₃ NPs are provided in this paper.

The use of tartaric acid as ligand in the BiFeO₃ synthesis produces nanopowders with crystallite size around 26 nm. The main byproduct formed is the metastable β -Bi₂O₃ phase, which could be efficiently removed with CH₃COOH. Also, some traces of tiny Bi₂Fe₄O₉ NPs of about 5 nm were found in the same sample. The thorough characterization of these BiFeO₃ NPs led us to the following relevant findings: (i) the presence of both types of electronic transitions (direct and indirect, located in the visible region) and (ii) a remarkable selective enhancement in the intensity of the BiFeO₃ A₁ mode (216 cm⁻¹), as a consequence of the resonance Raman effect. Both i and ii were observed and discussed for the first time in this work.

When glycine is used as ligand, a porous BiFeO₃ network, composed of tightly assembled and sintered nanocrystallites of high purity, was obtained. The following two important issues have not been addressed before in the literature: (i) There is an intermediate produced during the synthesis that could be the key species to generate high purity BiFeO₃. This intermediate is a porous network of nanoparticles smaller than 10 nm, mainly formed by BiFeO₃, an amorphous phase, and probably a nonstoichiometric BiFeO_{3-x}. This porous array formation is explained on the basis of the chemical reaction that involves an explosive glycine—nitrates mixture. (ii) There is a clear correlation between the singular morphology of the BiFeO₃—NPs—PNt and its magnetic properties, which are significantly different from those of its analogue derived from tartaric acid.

The BiFeO₃ nanoparticles, prepared using tartaric acid, have lower resonance field (and higher g-factor) than that of the BiFeO₃–NPs–PNt; this is the first report of the main EPR parameters for high purity BiFeO₃ nanopowders.

ASSOCIATED CONTENT

S Supporting Information

Additional figures and tables. Video of HRTEM images showing the interaction of the electron beam with 5 nm $\rm Bi_2Fe_4O_9$ NPs. This material is available free of charge via the Internet at http://pubs.acs.org.

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Notes

The authors declare no competing financial interest.

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